

Class H Statistical Mechanics

H1. Combinatorics Video Homework Pascal Triangle

H2. The Statistical Problem  $N$  particles  $n_1$  in energy  $\epsilon_1$   
 $n_2$  " "  $\epsilon_2$

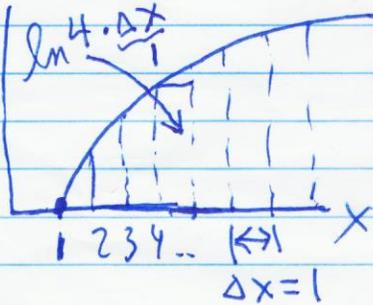
Maximize  $\Omega = \frac{N!}{n_1! n_2! n_3! \dots}$   
 # ways

By maximizing  $\ln \Omega$  Easier  $\rightarrow$   
 $N = n_1 + n_2 + n_3 + \dots$   
 $E = n_1 \epsilon_1 + n_2 \epsilon_2 + n_3 \epsilon_3 + \dots$

Maximize  $\Omega$  with the constraints  $N$  fixed and  $E$  fixed  $\hookrightarrow$  total energy

$\rightarrow$  We will maximize  $\ln \Omega = \ln N! - \ln n_1! - \ln n_2! - \dots$

$\ln x$   $\ln n! = \ln 1 + \ln 2 + \dots + \ln n$  In general  $N$  is super large  $n_i$  are large



$\ln n! \approx \int_1^n \ln x dx$

Integrate by parts  $d(x \ln x) = \ln x dx + \frac{1}{x} x dx$

$\ln n! = \int_1^n d(x \ln x) - \int_1^n dx = n \ln n - (n-1)$

Note  $\ln 1 = 0$   $x \ln x \Big|_1^n - x \Big|_1^n$   $\ln n! \approx n \ln n - n$   
 $n \ln n - \ln 1 - (n-1)$  Throw away the 1 for large  $n$

H3. Undetermined Multipliers Video Homework Approximation  $\ln n! \approx n \ln n - n$

H4. Maximize  $\Omega$   
 $\ln \Omega = N \ln N - N - n_1 \ln n_1 + n_1 + \dots$

$\ln \Omega \approx N \ln N - N - n_1 \ln n_1 - n_2 \ln n_2 - \dots$  So can't do  $\frac{\partial \Omega}{\partial n_i} = 0$  etc.  
 But  $dN = dn_1 + dn_2 + \dots = 0$   
 $dE = \epsilon_1 dn_1 + \epsilon_2 dn_2 + \dots = 0$

So we set  $d(\ln \Omega) - \alpha dN - \beta dE = 0$  Since  $dn_i$  are not all independent  
 The constants  $\alpha + \beta$  will enable us to consider  $dn_i$  independent

$-N$  cancels with  $n_1 + n_2 + \dots$

$\sum_i \left[ \frac{\partial \ln \Omega}{\partial n_i} - \alpha \frac{\partial N}{\partial n_i} - \beta \frac{\partial E}{\partial n_i} \right] dn_i = 0$   
 Can consider independent since  $\alpha + \beta$  are set to make it work  
 $\frac{\partial \ln \Omega}{\partial n_1} = \ln N - \ln n_1$

Macrostates have variables  $P, V, T, S$   
Microstates # ways to place particles in energy levels

$S = k \ln \Omega$   
(Shown on this page)  
Connects Macrostate to Microstates

$\frac{\partial \ln \Omega}{\partial n_i} = \ln N - \ln n_i$  True for each  $n_i$   $\uparrow$   $i^{th}$  case

$\ln N - \ln n_i - \alpha - \beta \epsilon_i = 0$

$\ln \frac{N}{n_i} = \alpha + \beta \epsilon_i$        $\ln \frac{n_i}{N} = -\alpha - \beta \epsilon_i$

$\frac{n_i}{N} = e^{-\alpha - \beta \epsilon_i}$

H5. Evaluating  $\alpha$  and  $\beta$

$n_i = N e^{-\alpha - \beta \epsilon_i}$

Find  $\alpha$   $N = \sum_i n_i = e^{-\alpha} N \sum_i e^{-\beta \epsilon_i} \Rightarrow e^{-\alpha} = \frac{1}{\sum_i e^{-\beta \epsilon_i}}$

Partition function  $Z = \sum_i e^{-\beta \epsilon_i}$

$n_i = \frac{N e^{-\beta \epsilon_i}}{Z}$  The  $\alpha$  is gone.

Find  $\beta$  Video Homework  $\rightarrow \beta = \frac{1}{KT}$   
Pascal's Law is used

H6. Entropy  $\Delta U = \Delta Q - \Delta W$  is the 1<sup>st</sup> Law of Thermodynamics  
 $\Delta Q = \Delta U + P \Delta V$   
 $\Delta Q = \frac{3}{2} nR \Delta T + \frac{nRT}{V} \Delta V$   
 $\Delta W \Rightarrow W$  doesn't. It depends on path.  
 $P, V, T$  have meaning for the state of the gas  
All good variables on right side

Divide by T makes good

$\frac{\Delta Q}{T} = \frac{3}{2} nR \frac{\Delta T}{T} + \frac{nR}{V} \Delta V$

$T \Delta S' = \Delta Q$  Entropy  $S'$  good macroscopic variable.

$dU = T dS' - P dV$

$\frac{\partial U}{\partial S'} = T$

But  $d(\ln \Omega) - \alpha dN - \beta dE = 0$

from before

$dE = \frac{1}{\beta} d(\ln \Omega) - \frac{\alpha}{\beta} dN$

$\frac{\partial E}{\partial \ln \Omega} = \frac{1}{\beta} = kT$

But  $U = E$

Therefore  $S' = k \ln \Omega$

Disorder # ways multiply

Since logs add

Note Entropies add combining 2 systems  $\Omega = \Omega_1 \cdot \Omega_2$